

Cavity Ringdown Spectroscopy of *cis-cis* HOONO and the HOONO/HONO₂ Branching Ratio in the Reaction OH + NO₂ + M

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The termolecular association reaction OH + NO₂ + M was studied in a low-pressure discharge flow reactor, and both HONO₂ and HOONO products were detected by infrared cavity ringdown spectroscopy (IR-CRDS). The absorption spectrum of the fundamental ν_1 band of the *cis-cis* isomer of HOONO (pernitrous or peroxyxynitrous acid) was observed at 3306 cm⁻¹, in good agreement with matrix isolation studies and ab initio predictions. The rotational contour of this band was partially resolved at 1 cm⁻¹ resolution and matched the profile predicted by ab initio calculations. The integrated absorbances of the ν_1 bands of the *cis-cis* HOONO and HONO₂ products were measured as a function of temperature and pressure. These were converted to product branching ratios by scaling the experimentally observed absorbances with ab initio integrated cross sections for HOONO and HONO₂ computed at the CCSD(T)/cc-pVTZ level. The product branching ratio for *cis-cis* HOONO to HONO₂ was $0.075 \pm 0.020(2\sigma)$ at room temperature in a 20 Torr mixture of He/Ar/N₂ buffer gas. The largest contribution to the uncertainty is from the ab initio ratio of the absorption cross sections, computed in the double harmonic approximation, which is estimated to be accurate to within 20%. The branching ratio decreased slightly with temperature over the range 270 to 360 K at 20 Torr. Although *trans-perp* HOONO was not observed, its energy was computed at the CCSD(T)/cc-pVTZ level to be $E_0 = +3.4$ kcal/mol relative to the *cis-cis* isomer. Statistical rate calculations showed that the conformers of HOONO should reach equilibrium on the time scale of this experiment. These results suggested that essentially all isomers had converted to *cis-cis* HOONO; thus, the reported branching ratio is a lower bound for and may represent the entire HOONO yield.

Introduction

The free radical chain termination reaction



is important in oxidation chemistry throughout the lower atmosphere. Radicals in both the HO_x and NO_x families mediate the destruction of ozone in the stratosphere and are central to the formation of photochemical smog in the troposphere.^{1,2} Reaction 1 reduces the concentrations of these species by converting them into the stable reservoir species nitric acid. The

rate of reaction 1 has important consequences for tropospheric ozone production and photochemical smog. In polluted urban airsheds with high ratios of NO_x to volatile organic compounds (VOC), the rate of reaction 1 is one of the single most influential parameters determining peak ozone mixing ratios.^{3,4} At high enough NO_x levels, urban ozone levels decrease with increasing NO_x, because elevated NO_x levels reduce OH concentrations through reaction 1 and lead to reduced VOC oxidation rates. Tonnesen⁵ has shown that model predictions of ozone production can be almost inversely proportional to changes in the rate coefficient for reaction 1 under these conditions. This negative feedback complicates ozone attainment strategies, which must often balance regulation of VOC vs NO_x emissions. Furthermore, the uncertainty in the rate of reaction 1 is the major contribution to model uncertainties in predicted ozone at high NO_x levels. This uncertainty propagates to scales such as the Maximum Incremental Reactivity, a measure of the ozone-

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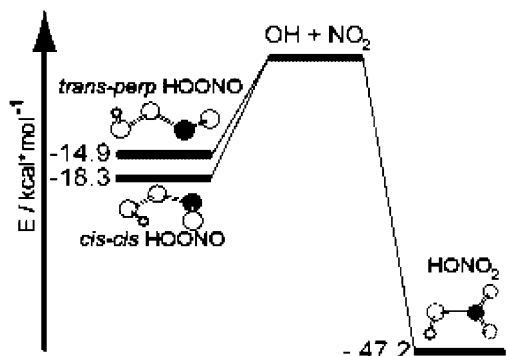


Figure 1. Energy level diagram of the reaction $\text{OH} + \text{NO}_2$. The *cis-cis* and *trans-perp* isomers of HOONO are shown. Energies for HOONO are from CCSD(T) calculations.

forming potential of individual VOCs.^{6,7} In the stratosphere, both HO_x and NO_x radicals are involved in catalytic cycles that deplete ozone. Reaction 1 couples these catalytic cycles, and is one of several key reactions that lead to a more complex dependence of stratospheric ozone on NO_x levels. Sensitivity analyses have shown that differences in recommended rates for reaction 1 have a significant effect on mid-latitude ozone concentrations.⁸

Despite extensive studies over a wide range of temperatures and pressures,^{9–26} significant uncertainties remain at atmospherically relevant pressures and temperatures. Much attention has recently focused on the yield for a second minor product channel



producing HOONO (peroxynitrous or pernitrous acid), a weakly bound isomer of nitric acid. The NASA Panel for Data Evaluation²⁷ and the IUPAC Gas Kinetic Data Evaluation^{28,29} summarize the current evidence for this channel, but neither source currently makes a recommendation for its inclusion in atmospheric models. A significant branching fraction into channel 2 is of importance in the atmosphere. Formation of HOONO would reduce the efficiency of the $\text{OH} + \text{NO}_2 + \text{M}$ reaction as a sink for OH and NO_2 radicals, because HOONO would rapidly dissociate or photolyze back to reactants.

While Leighton³⁰ had suggested the existence of the HOONO channel in the gas phase, Robertshaw and Smith⁹ first proposed that a significant HOONO yield could resolve a discrepancy in the pressure dependence of reaction 1. The effective bimolecular rate constant k_1 has an anomalous falloff curve; the observed high-pressure rate constant k_1^∞ is significantly larger than the high-pressure limit extrapolated from low-pressure data. The discrepancy can be removed if the HOONO channel is insignificant at low pressures, but becomes comparable to the HONO_2 channel at high pressures. Golden and Smith³¹ modeled the $\text{OH} + \text{NO}_2$ reaction and showed that the high- and low-pressure data could be described quantitatively by statistical models that included HOONO. Hippler et al.³² have recently revised the high-pressure rate constants in He and N_2 downward by $\approx 30\%$, but the discrepancy remains.

An energy level diagram illustrating these two channels is shown in Figure 1. At low pressures, the initial HOONO association adduct will redissociate more quickly than HONO_2 due to the former's lower density of states, and bound HOONO will be formed in low yields. At high pressures, the weakly bound complex will be stabilized rapidly and the HOONO yield could be comparable to, or even exceed, the HONO_2 yield.

Golden and Smith,³¹ Matheu and Green,³³ and Troe³⁴ have modeled the pressure and temperature dependence of k_1 and k_2 . Since the submission of this paper, there has been a new set of master equation calculations by Golden, Barker, and Lohr.³⁵ While all sets of calculations suggest that HOONO formation should compete with the HONO_2 channel in the reaction $\text{OH} + \text{NO}_2$, the resulting branching ratios at 1 atm and 298 K vary from $> 1\%$ (Matheu and Green), to 2.5% (Troe), to 30% (Golden and Smith). However, rate constants, falloff curves, and product yields cannot be calculated quantitatively from first principles, because our knowledge of the transition states, relaxation processes, and potential surfaces (especially for HOONO formation) is too limited. These models rely to some extent on fitting experimental data, and thus may not accurately predict branching ratios without direct measurements of HOONO.

There have been extensive calculations^{36–45} of the energetics and conformations of peroxynitrous acid. From MP2 calculations, McGrath and Rowland found three stable forms of HOONO, the *cis-cis*, *cis-trans*, and *trans-perp* isomers (distinguished by the HOON and OONO torsional angles). However, recent results at the highest levels of theory (QCISD(T)/cc-pVTZ) indicate that only the *cis-cis* and *trans-perp* conformations are stable stationary points.⁴⁶ The best purely theoretical estimate of the O–O bond energy of *cis-cis* HOONO is $\Delta_2 H_0 = -19.8$ kcal/mol at the CCSD(T)/CBS limit with corrections.⁴⁴ This estimate is based on the computed $\Delta_f H_0^0(\text{NO}_2)$, which is in error by 1.5 kcal/mol. A perhaps more accurate assessment, based on $\Delta_f H_0^0(\text{cis-cis HOONO}) - \Delta_f H_0^0(\text{HONO}_2) = 29.0$ kcal/mol, gives $\Delta_2 H_0 = -18.3$ kcal/mol. In either case, this binding energy is large enough to expect that *cis-cis* HOONO will be a relatively stable molecule at room temperature (see Figure 1). The energy of the *trans-perp* isomer is computed to be 1–3.2 kcal/mol above the *cis-cis* form. The transition state for isomerization between these isomers, the *perp-perp* isomer, has an MP2 energy (including ZPE correction) of 11.4 kcal/mol relative to the *cis-cis* isomer.³⁷

HOONO has been studied extensively in the condensed phase, though the most direct spectroscopic studies come from matrix-isolation work. Lee and colleagues^{47–50} reported vibrational and UV absorption spectra of HOONO trapped in rare gas matrices. The observed infrared bands matched frequencies and intensities predicted by ab initio calculations for the *cis-cis* and *trans-perp* isomers. The *cis-perp* isomer is predicted to have vibrational frequencies similar to those of the *trans-perp* isomer, but it was ruled out on the basis of intensity patterns.

In the gas phase, two earlier attempts at infrared absorption spectroscopy failed to detect HOONO. Burkholder et al.⁵¹ attempted to prepare HOONO in a fast-flow multipass cell using a variety of methods ranging from microwave discharge to direct synthesis with H_2O_2 and NOBF_4 . They estimated an upper limit for HOONO yield of 5% based on sensitivity arguments, and an HONO_2 yield of $75\% + 25\%/-10\%$ from mass balance. More recently, Dransfield et al.⁵² attempted to detect HOONO formed by reaction 1 with reaction modulation FTIR spectroscopy, but found no unidentified bands that could be assigned to HOONO in the range $700\text{--}3200$ cm^{-1} . They estimated an upper limit for HOONO of 3–4% based on mass balance arguments.

Direct evidence for formation of HOONO in the gas phase by reaction 2 has been reported only recently. In 2001, Donahue et al.⁵³ described a series of ^{18}OH isotopic scrambling experiments that supported the existence of reaction 2 and led to an estimate of the low-pressure branching ratio of $k_2^\circ/k_1^\circ = 0.17$ at room temperature. Nizkorodov and Wennberg⁵⁴ then reported

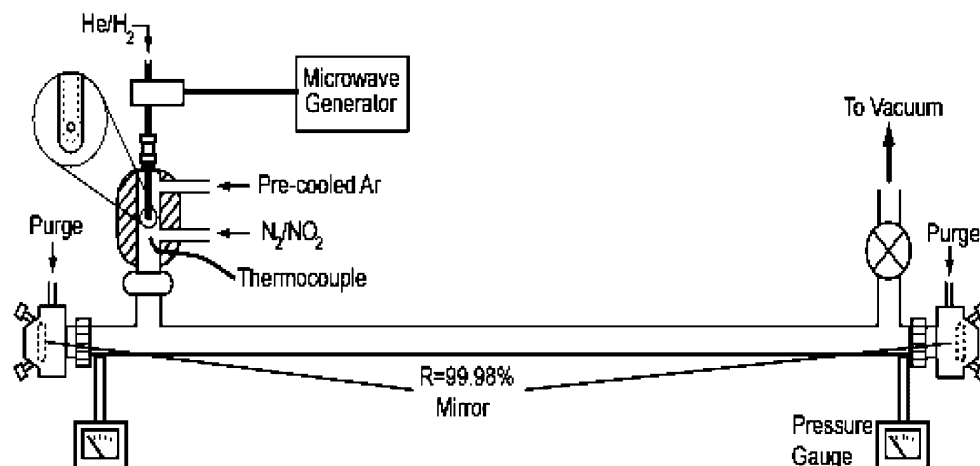


Figure 2. Overview of the experimental apparatus. A microwave discharge is used to produce H atoms that are introduced into an Ar flow through the injector tip (exploded view), which then mixes with an NO_2/N_2 flow. The H atoms react by $\text{H} + \text{NO}_2 \rightarrow \text{OH} + \text{NO}$ to form OH radicals. The OH radicals react with NO_2 to form HOONO and HONO₂, within 1 cm. Infrared absorption spectra of products are detected downstream in the ringdown cavity, formed by two high-reflectivity mirrors (dashed lines, $R = 99.98\%$ at $2.8 \mu\text{m}$).

the first spectroscopic observation of gas-phase HOONO from the reaction of OH and NO_2 . They observed the $2\nu_1$ overtone spectrum of HOONO formed in a low-pressure discharge flow cell using action spectroscopy, detecting OH by LIF after predissociation of the overtone-excited molecules. They observed several bands which they tentatively assigned to the isomers of HOONO and combination bands. They estimated that the HOONO yield was $5 \pm 3\%$ at 253 K and 20 Torr. In 2002, Hippler et al.³² reported kinetic evidence for channel 2. They observed a biexponential decay in the OH LIF signal in the reaction $\text{OH} + \text{NO}_2$ at 5–100 bar and 430–475 K, which they modeled as rapid loss of OH by reactions 1 and 2, followed by slower decomposition of HOONO by the reverse rate k_{-2} . They found the enthalpy of reaction 2 at this temperature range to be $\Delta H_2(0) = -19.8 \text{ kcal/mol}$, in good agreement with the theoretical predictions of the *cis-cis* HOONO dissociation energy D_0 .

In principle, the branching ratio between reactions 1 and 2 can be measured by infrared absorption spectroscopy of both HOONO and HONO₂, but the difficulty of detecting HOONO at such low concentrations ($<10^{12} \text{ cm}^{-3}$) has been compounded by possible spectral interference from the high concentrations of HONO₂ and NO_2 . However, the lowest energy isomer, *cis-cis* HOONO, can be distinguished from HONO₂ in the $3 \mu\text{m}$ region, because the OH stretch is significantly red-shifted relative to that of HONO₂. Ab initio calculations indicate that this shift arises from a hydrogen-bonding interaction between the H atom and the terminal O atom which occurs only in the *cis-cis* conformation (see Figure 1).^{36,37} Lo and Lee observed this band at 3285 cm^{-1} in Ar matrix.⁵⁵ In contrast, they observed the ν_1 band of the *trans-perp* isomer at 3563.4 and 3545.7 cm^{-1} (different substitution locations), nearly coincident with the OH stretch of nitric acid at 3550.0 cm^{-1} .

In this work, we report a joint experimental/theoretical study to measure the *cis-cis* HOONO-to-HONO₂ branching ratio from the reaction $\text{OH} + \text{NO}_2 + \text{M}$, which we denote $k_2(c-c)/k_1$, in a low-pressure discharge flow reactor. We have exploited the high sensitivity of infrared cavity ringdown spectroscopy (IR-CRDS) to observe directly the absorption spectra of *cis-cis* HOONO as well as HONO₂ in the $3 \mu\text{m}$ region. The branching ratio can be determined from the integrated intensities of the OH stretch modes of both species with the aid of ab initio vibrational band strengths calculated at the QCISD and CCSD(T)/cc-pVTZ level. Finally, we perform statistical calculations

to predict the rate that the *trans-perp* HOONO isomer converts to the *cis-cis* form and compute the *trans-perp* HOONO energy. From these results, we infer that the observed *cis-cis* yield most likely represents the total HOONO yield.

Experiment

An overview of the experimental apparatus is shown in Figure 2. The flow cell comprised two regions, a discharge flow reactor in which the $\text{OH} + \text{NO}_2$ reaction took place, and a long path spectroscopic detection region. In the reaction region, hydrogen atoms were formed by microwave discharge and thermalized by injection into a flow of pre-cooled argon. This gas was then mixed with NO_2 seeded in N_2 to produce OH radicals by the reaction



with NO_2 in excess. The OH products of reaction 3 then reacted with the NO_2 . The resulting mixture of products flowed into the detection region, where they were detected by infrared absorption from 2.6 to $3.4 \mu\text{m}$ using cavity ringdown spectroscopy. Experiments were performed over the pressure range 5 to 40 Torr and the temperature range 240 to 360 K.

The cell was constructed of Pyrex and quartz, and the interior walls were coated with halocarbon wax to minimize wall reactions. The entire system was evacuated by a 52-cfm mechanical pump through a stopcock and an LN_2 -trapped foreline. Two capacitance pressure transducers were located at either end of the detection region to verify constant pressure throughout the cell during measurements. All of the gases used in these experiments were UHP grade (99.999%) or higher purity. Six mass flow transducers coupled with needle valves enabled us to vary the flows of the reactant gases as well as the purge flows over the ringdown mirrors. The flows and estimated concentrations of reactants are listed in Table 1. The linear flow velocity was typically 100–200 cm/s, and the total residence time in the cell was typically 200 ms.

A. OH + NO_2 Reaction Region. H-atoms were produced in a microwave discharge source. H_2 was mixed in a 1:200 ratio with He that had been passed through a LN_2 -cooled zeolite trap to remove condensable impurities. The mixture then flowed through a quartz capillary (4-mm i.d.) in which a discharge was maintained by an Evenson microwave cavity. The capillary extended collinearly into the main reactor cell (2.5-cm i.d.) in

TABLE 1: Typical Flows and Concentrations for Branching Ratio Measurements

species	flow/sccm	concentration/ cm ⁻³ (300 K, 20 Torr)
He	615	2.85×10^{17}
Ar	515	2.38×10^{17}
N ₂	242	1.12×10^{17}
NO ₂	13	6.0×10^{15}
H ₂	3	1.38×10^{15}

which pre-cooled Ar gas flowed. The discharge gas was injected in jets perpendicular to the Ar flow through small holes near the capillary tip. Turbulent mixing of the discharge gas with the Ar flow ensured an even distribution of H atoms in the reactor and helped to dissipate heat from the discharge.

A mixture of 4% NO₂ in N₂ by volume (Matheson, certified gravimetrically) was introduced through a 4-mm i.d. inlet approximately 1 cm downstream from the end of the quartz discharge tube. We adjusted the location of the H-atom injection with respect to the N₂/NO₂ entrance to allow for the mixing and cooling of the discharge flow, while keeping wall loss of H-atoms to a minimum. At the concentrations used in this experiment, the reaction of H and NO₂ was essentially complete within microseconds, while the subsequent association reaction of OH + NO₂ + M was complete in <500 μs. Thus, products from reactions 1 and 2 were formed within a few millimeters along the flow, well before the gas entered the spectroscopy region. Although H + NO₂ is known to produce high yields of vibrationally excited OH,⁵⁶ NO₂ is an effective quencher. Furthermore, OH(*v*) + NO₂ will produce highly excited transient adducts which will not stabilize to form products.

The temperature of the reaction region was measured with a calibrated type “T” thermocouple placed in the center of the reaction zone. The temperature at the point of reaction was slightly higher than the reactor walls, because the He/H₂/Ar mixture had not fully equilibrated. Thermal gradients led to temperature variations of up to ±3° C across the reaction region; these uncertainties are reflected in abscissa error bars in the temperature dependence.

B. Detection of HOONO and HONO₂. Infrared absorption spectra of the reaction products were obtained in a long-path cell situated downstream of the reaction region, 8 cm from the NO₂ injection point. Most spectra were recorded in a room-temperature cell (1.9-cm i.d.) which had a cavity length $L = 81.5$ cm and a sample absorption path length $L_s = 51$ cm through the reactant gas mixture. The absorption path length L_s was determined by the distance from the gas inlet port to the pump-out port. We occasionally used a second jacketed cell ($L = 67$ cm, $L_s = 32$ cm, 2.5-cm i.d., 20 cm distance from the NO₂ inlet port), which allowed us to vary the detection cell temperature in order to investigate temperature-dependent wall loss and residence time effects.

Absorption spectra of the reaction products were recorded by pulsed cavity ringdown spectroscopy.^{57–59} A mid-infrared laser pulse was injected into an optical cavity made by two high-reflectivity infrared mirrors ($R = 99.98\%$, centered at $\lambda = 2.8$ μm, 6 m radius of curvature) sealed to the ends of the detection cell. Each mirror was separated from the main flow of gas by a purge volume approximately 15 cm long. Dry argon flowed over the mirrors (8 sccm over each mirror) and through the purge volumes to protect the mirrors and to minimize losses in reflectivity due to deposition.

We produced tunable (2.6–3.4 μm) mid-IR in a two-stage, optical parametric amplifier (OPA)⁶⁰ by mixing the second harmonic of an Nd:YAG laser and the output of a tunable dye

laser in a pair of KTP crystals.⁶⁰ The OPA produced an IR beam with energy $E \approx 1$ mJ/pulse and diameter $d \approx 2$ mm after transmission through three filters to remove the residual 532 and 630 nm light. The IR line width was $\Gamma \approx 1$ cm⁻¹. The IR wavelength was typically calibrated using the frequencies of known HONO and HNO₃ peaks in the spectrum. The IR beam was aligned into the cavity, and the residual beam exiting the cell was focused tightly onto a 1-mm diameter InSb detector (Infrared Associates). Each ringdown trace was amplified and then accumulated with a 300 MHz bandwidth digital oscilloscope with 8-bit digitizer/12-bit averager (Tektronix TDS 320). For each point in a spectrum, typically 32 ringdown traces were averaged by the oscilloscope before transfer to the PC, which then fit the data to a single-exponential function and constant baseline by the Levenberg–Marquardt algorithm.⁶¹

The ringdown lifetime of radiation trapped in an evacuated optical cavity of length L with mirrors of reflectivity R is $\tau_0 = L/c(1 - R)$, where c is the speed of light.^{62–64} Typical empty-cell ring-down times were $\tau_0 \approx 10$ μs, in good agreement with the stated mirror reflectivity (99.98% at the peak wavelength). In the presence of an absorber, the lifetime is given by $\tau = L/c(1 - R + \alpha L_s)$, where $\alpha = [X]\sigma_X$, σ_X is the extinction coefficient of the absorber X, and L_s is the absorption path length through the sample. The measured extinction coefficient is then

$$\alpha = \frac{L}{cL_s} \left(\frac{1}{\tau} - \frac{1}{\tau_0} \right)$$

The possible dependence of the measured extinction coefficient on the concentration of absorber will be discussed in detail below.

Spectra of reaction products were typically recorded by scanning the spectral region of interest with the discharge on and off, and subtracting the two scans. This subtraction eliminated background due to species present in the absence of reaction, as well as to the wavelength-dependent ringdown loss of the mirrors.

We estimated the sensitivity of our CRDS setup from $\Delta\tau$, the minimum detectable change in τ .⁶³ In a typical run, we found $\Delta\tau/\tau_0 \approx 2\%$ (averaging 32 laser shots), giving a minimum detectable absorption of $\alpha_{\min} \approx 1 \times 10^{-7}$ cm⁻¹ (1.8×10^{-7} cm⁻¹ √Hz) for mirrors with reflectivity $R = 99.97\%$. For the detection of HONO₂ with the ν_1 band, which has an integrated intensity of 9.5×10^{-18} cm/molecule,⁶⁵ the minimum detectable concentration of nitric acid was $[\text{HONO}_2]_{\min} \approx 5 \times 10^{11}$ molecules/cm³ at a single point in the R-branch. Integration over the OH stretching band further improved the sensitivity by roughly $N^{1/2}$, where $N \approx 20$ is the number of data points making up fwhm of the band.

We investigated the effect of the purge gas flow rate on the detection of products. The linear flow velocity of Ar through each of the purged volumes was ≈ 3 cm/s, or 0.6% of the total flow. We measured both product absorption signals while independently varying the purge flow at the inlet and outlet ends of the detection cell from 0.5% to 9%. There was a gradual decrease (approximately 20%) in the absorption by both HONO₂ and HOONO as we increased the purge flow rates to 9%, but the fractional yield of HOONO remained constant. Thus, while some reaction gas mixture entered the purged volume at the flow rates used in this experiment, these gases were flushed out quickly and did not affect branching ratio measurements.

Computational Methodology

A. Ab Initio Calculations of Structure, Energetics, and Intensities. Ab initio electronic structure calculations were

TABLE 2: Computed Molecular Constants for the *cis-cis* and *trans-perp* Conformational Isomers of HOONO

	cis-cis				trans-perp		
	QCISD cc-pVTZ ^a	QCISD(T) cc-pVTZ ^a	CCSD(T) cc-pVTZ	MP2 6-31G(d) ^b	QCISD cc-pVTZ	CCSD(T) cc-pVTZ	MP2 6-31G(d) ^b
ΔE_0 (kcal/mol)						3.4	2.9 ^c
A (MHz)	21726	21262	21310	20900	55716	54127	52576
B	8133	8041	8081	7949	5060	4977	4902
C	5917	5835	5859	5759	4709	4627	4550
ν_1 (cm ⁻¹)		3496(a')	3521	3547 ^d		3780	3841 ^e
ν_2		1649(a')	1630	1595		1736	1720
ν_3		1461(a')	1458	1476		1412	1408
ν_4		939(a')	969	991		992	999
ν_5		796(a')	838	858		816	811
ν_6		618(a')	723	776		497	443
ν_7		397(a')	419	418		366	373
ν_8		515(a'')	523	532		302	311
ν_9		363(a'')	383	360		211	219

^a Refs 43,44. ^b McGrath and Rowland, ref 37. ^c MP2/6-311+G(3df,2p), ref 37. ^d Frequencies calculated with the 6-311G(2df,2p) basis set. ^e Frequencies calculated with the 6-311G(d,p) basis set.

performed to predict (a) structures, vibrational frequencies, and energies of the *cis-cis* and *trans-perp* isomers of HOONO, and (b) band intensities of HOONO, HO₂NO₂, and HONO₂. Calculations were performed with the Gaussian 98⁶⁶ and MOLPRO⁶⁷ programs. Coupled cluster theory⁶⁸ with single and double excitations, with perturbative estimates of connected triples CCSD(T), was used along with the quadratic configuration interaction method (QCISD(T) and QCISD).⁶⁹ The QCISD(T) method is correct through fourth-order terms, and it includes the same fifth-order terms as the CCSD(T) method. In the absence of large multireference effects or large single excitation amplitudes, the QCISD(T) method should be very reliable,⁶⁹ and its performance is very similar to that of CCSD(T).⁷⁰ The calculations were performed with the Dunning correlation-consistent polarized valence basis sets. The two basis sets used throughout this study are the cc-pVDZ basis set (3s2p1d for C, N, and O; 2s1p for H) and then the cc-pVTZ basis set (4s3p2d1f for C, N, and O; 3s2p1d for H).^{71,72} From many benchmark calculations with extended basis sets, the CCSD(T) and QCISD(T) methods have shown that high accuracy can be achieved for the molecular structure and vibrational frequencies.^{73,74} Vibrational frequencies were calculated with both the QCISD and CCSD(T) methods. Intensities were calculated with the QCISD method and scaled to estimate the CCSD(T) intensities by the factor 0.968. The scale factor was determined from a calibration set of molecules of well-known intensities and also from calculation of intensities for the calibration set of molecules using the CCSD(T) method. The CCSD(T) intensities were calculated with MOLPRO. The calibration set was HCN, H₂O, CH₄, and CH₃O. A total of 14 intensities were used in the scale factor determination.

B. Transition Dipole Moment Vector. Additional calculations were carried out to determine the orientation of the vibrational transition dipole moment for the ν_1 band of HOONO. The electrostatic dipole moment was computed at the B3LYP/cc-pVTZ level using Gaussian 98 for several displacements from equilibrium along the Q_1 normal coordinate using the L matrix given at the equilibrium geometry. The resulting dipole vector functions were fit to a polynomial and then numerically integrated over harmonic oscillator wave functions to obtain the three components of the transition dipole matrix element in the principal axis frame.

C. Calculations of Isomerization Rates. We computed the formation and isomerization rates of the isomers of HOONO using the program MultiWell.^{75,76} We focused in particular on the rate of isomerization from the *trans-perp* isomer to the lower

energy *cis-cis* isomer. For the current calculation, we used parameters for the *perp-perp* transition state for rotation of the OONO dihedral angle, which lies between the *trans-perp* and *cis-perp* isomers. The energy of the *perp-perp* isomer was computed at the MP2/6-311G(3df,2p) level with MP2/6-31G-(d) frequencies by McGrath and Rowland (note that the MP2 energies listed in their Table 2 are not corrected for zero-point energy).³⁷ We employed the *trans-perp* and *cis-cis* vibrational frequencies from the same calculation, but we used our CCSD(T) calculation of the *trans-perp* energy. We also performed some calculations of the initial branching ratios for reactions 1 and 2, using transition state parameters based on those of Golden and Smith.³¹ Energy transfer collisions were described by the single-exponential model⁷⁵ with $\alpha = 45 \text{ cm}^{-1} + 0.005 E_{\text{upper}}(\text{cm}^{-1})$, and Lennard-Jones parameters $\epsilon = 180 \text{ K}$ and $\sigma = 4.2 \text{ \AA}^2$.

Results

A. Spectroscopy of OH + NO₂ Products. Figure 3 shows a room-temperature scan from 3200 to 3700 cm⁻¹ taken in 0.2 cm⁻¹ steps with the flow cell conditions shown in Table 1. The full-scale spectrum, shown in the inset, is dominated by the ν_1 HONO₂ transition. With the discharge off, the HONO₂ spectrum was approximately 8–10 times weaker. This small background HONO₂ signal, which was always present in the cell during the experiments, arose largely from slow desorption from passivated walls. Several other absorption bands were evident in this region. All of the bands had partially resolved rotational structure. The smaller bands could be assigned to known transitions: the $2\nu_2$ band of HONO₂, the $2\nu_2$ and ν_1 bands of both *cis*-HONO and *trans*-HONO, and the $2\nu_3$ transition of NO₂. There remained one unidentified band centered at 3306 cm⁻¹.

B. Assignment of the *cis-cis* HOONO Band. Figure 4 shows the spectrum from 3250 to 3375 cm⁻¹ after subtraction of the empty-cell background. The sharp features at 3260 and 3372 cm⁻¹ were assigned to the Q-branches of the $2\nu_2$ transitions of *cis*- and *trans*-HONO, respectively. A clearly defined set of P, Q, and R subbands centered at 3306 cm⁻¹ was identifiable between these features. We assigned this band to the $\nu_1 = 0 \rightarrow 1$ (O–H stretch) transition of *cis-cis* HOONO for the following reasons.

(1) Chemistry. This band was not observed if the discharge was off, even with all reagent gases (Table 1) present. We did not observe any absorption in this region when we recorded experimental reference spectra of either HONO or HONO₂

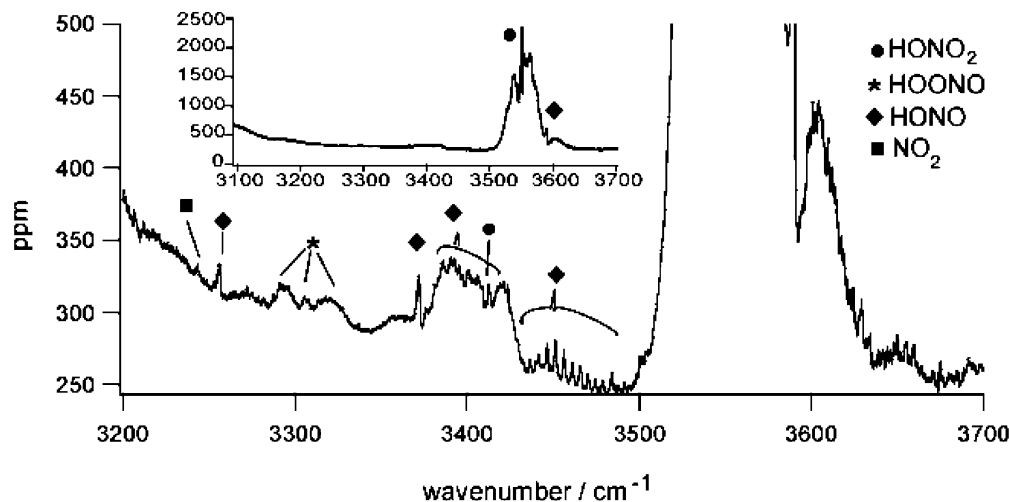


Figure 3. Absorption spectrum of the products of OH + NO₂ at 20 Torr and 300 K recorded by infrared cavity ringdown spectroscopy. No background has been subtracted; the varying baseline is due to the underlying variations in ringdown mirror reflectivity. The laser line width is 1 cm⁻¹ and the step size of the scan is 0.2 cm⁻¹. The full scale spectrum, shown in the inset, shows that nitric acid is the dominant product in the reaction OH + NO₂. Weaker bands in the magnified scan assigned to known transitions are marked: HONO₂ (●), HONO (◆), or NO₂ (■). The only unassigned band, at 3306 cm⁻¹, is attributed to the ν_1 transition of *cis-cis* HOONO (*).

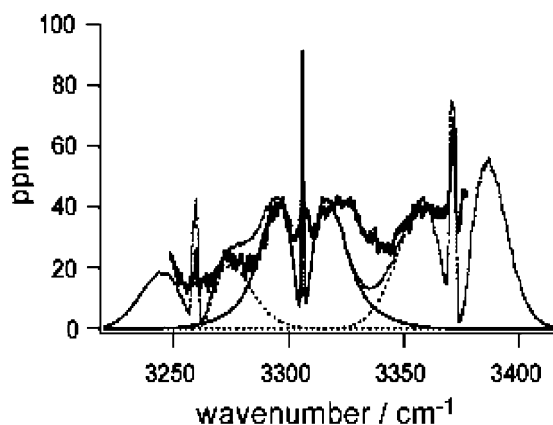


Figure 4. Spectrum of the *cis-cis* HOONO ν_1 band and adjacent $2\nu_2$ bands of *cis* and *trans*-HONO, with the ringdown baseline subtracted. The thick line is the experimental data, and the gray line is the sum of three simulated spectra: *cis-cis* HOONO (thin line, centered at 3306 cm⁻¹), *cis*-HONO (short dash, 3260 cm⁻¹), and *trans*-HONO (long dash, 3372 cm⁻¹).

generated from bubbled samples. Furthermore, the band was not present if either NO₂ or H₂ was removed with the discharge on. We found no detectable spectra of H₂O₂ or HO₂, which have band origins closer to 3400 cm⁻¹, or of HO₂NO₂, which has an absorption at 3540 cm⁻¹. Thus, this band arose from a reaction product but could not be attributed to any of the primary reactant or product species (HONO, NO, NO₂, HONO₂) or other conceivable products (H₂O₂, HO₂NO₂).

(2) Deuterium Substitution. When D₂ was substituted for H₂ in the discharge flow, the band at 3306 cm⁻¹ disappeared but the two adjacent DONO $2\nu_2$ (NO₂ stretch overtone bands) were still readily observable. These results indicate that this mode most likely involved a hydride X–H stretch.

(3) Band Position and Intensity. The frequency of the band origin, 3306 cm⁻¹, red-shifted relative to the OH stretch of HONO₂ by 245 cm⁻¹, was in good agreement with the position expected for the *cis-cis* isomer of HOONO from both experiment and theory. The observed band center was close to the 3285 cm⁻¹ origin of the ν_1 fundamental of *cis-cis* HOONO observed in matrix isolation experiments; the difference can be accounted for by matrix-induced shifts. It also matched the fundamental position estimated by Nizkorodov and Wennberg,⁵⁴

based on their preliminary assignment of a band at 6365 cm⁻¹ to the $\nu_1 = 0 \rightarrow 2$ overtone transition of *cis-cis* HOONO. Their overtone assignment, taken in conjunction with our observation of the fundamental band, would give $\omega_e = 3430$ cm⁻¹ and $\omega_e x_e = 62$ cm⁻¹ for the ν_1 band of HOONO. This harmonic frequency lies close to $\omega_e = 3521$ cm⁻¹, predicted by our CCSD(T)/cc-pVTZ calculations for the *cis-cis* isomer of HOONO. The observed frequency did not match the OH stretch frequencies of the other stable isomer(s) (or of the predicted saddlepoints) of HOONO, all of which have frequencies close to that of nitric acid ($\omega_{\text{exp}} = 3550$ cm⁻¹).

Finally, the integrated band intensity was approximately 3% of the integrated intensity of the ν_1 band of HNO₃. This was consistent with expectations of a relatively low HOONO yield, assuming the two bands had similar absorption cross sections. A quantitative assessment of the HOONO yield is given below.

C. Ab Initio Calculations. There have been a number of ab initio calculations on the isomers of HOONO. Here we performed QCISD, QCISD(T), and CCSD(T)/cc-pVTZ calculations on *cis-cis* and *trans-perp* conformers of HOONO, to have a consistent set of high-level spectroscopic constants for the two bound isomers. A third isomer, the *cis-perp* conformation, was not found to be stable. Computational results on the torsional potentials will be reported elsewhere.⁴⁶

Table 2 lists results for the rotational constants and vibrational frequencies of the two isomers, along with a comparison with the MP2 calculations of McGrath and Rowland.³⁷ The numbering of the normal modes of the *trans-perp* isomer are ordered in descending frequency, while the modes of the *cis-cis* isomer, which has *C_s* symmetry, are based first on the symmetry *a'* and then *a''*. In general, the frequencies and rotational constants are close to the MP2 results.

In addition, we report the CCSD(T)/cc-pVTZ energy of the *trans-perp* isomer relative to the lowest energy *cis-cis* geometry to be $\Delta E_0 = +3.4$ kcal/mol, including zero-point correction. The dissociation energy of the *cis-cis* isomer has been reported at a higher level of theory; the best estimate is 18.3 kcal/mol.⁴⁴

To determine the relative yields of HOONO and HONO₂, we needed the relative absorption cross sections of the respective fundamental ν_1 transitions. The band strength of nitric acid has been measured, but not that of HOONO. We therefore relied on ab initio intensities computed at equivalent levels of theory

TABLE 3: Calculated Integrated Intensities for ν_1 (O–H Stretch) Vibrational Transitions in km mol^{-1} , and Comparison with Previous Experimental Values

	HONO ₂	cis-cis HOONO	trans-perp HOONO	HOONO ₂
QCISD/cc-pVTZ ^a	98.6	34.3	50.1	52.8
CCSD(T)/cc-pVTZ	95.4	33.2	48.6	51.2
experiment		57.0 ± 2.8^b		36.2 ± 4.4^c

^a Scaled by 0.968 to estimate CCSD(T) intensities (see text). ^b Ref 65. ^c Ref 87.

to determine the ratio of the absorption strengths of *cis-cis* HOONO and HONO₂. To estimate the error that may be associated with such a calculation, we also computed the intensity for pernitric acid, HO₂NO₂, a molecule of similar composition which, like HOONO, possesses an OH stretch coupled to a peroxide bond.

Table 3 shows the integrated intensities for the OH stretches in HOONO, HONO₂, and HO₂NO₂ computed with the QCISD(T)/cc-pVTZ and CCSD(T)/cc-pVTZ methods, as well as a comparison to published values. The QCISD(T) calculations were scaled to estimate the CCSD(T) intensities. While the calculated intensities of HONO₂ and HO₂NO₂ are greater than the experimental values by 67% and 41%, respectively, the computed HONO₂-to-HO₂NO₂ intensity ratios are 1.86 at the CCSD(T) level of theory, while the experimental ratio is 1.57. Thus, the computed ratios—though not the absolute magnitudes—agree with experiment to within 18%.

The CCSD(T) calculations gave an intensity ratio of 2.87 for HONO₂ to *cis-cis* HOONO, and 1.96 for HONO₂ to *trans-perp* HOONO; the scaled QCISD(T) ratios were nearly identical. The integrated absorbance of HOONO must be corrected by this ratio to determine the relative HOONO-to-HONO₂ product concentrations. The comparison with pernitric acid suggests that this value should be accurate to within roughly 20%. This estimation of the error is based on a single calibration point; hence, we cannot estimate the statistical uncertainty, but IR absorbances calculated at this level of theory, even within the double harmonic oscillator approximation, are typically accurate to within 20%.^{77,78} By estimating the ratio of intensities of similar stretch modes (OH stretches), systematic errors such as anharmonicity are at least partially corrected, though the *cis-cis* HOONO has an apparent hydrogen-bonding interaction not present in either HONO₂ or HO₂NO₂.

Lower-level B3LYP calculations were performed to predict the orientation of the vibrational transition dipole moment for the $\nu_1 = 0 \rightarrow 1$ transition of *cis-cis* HOONO. The transition dipole moment vector computed as described above was $(\mu_a, \mu_b, \mu_c) = (0.041 \text{ D}, 0.064 \text{ D}, 0.0 \text{ D})$. This vector differs slightly from the direction of the OH bond. This calculation then predicts a hybrid band with a 2.3:1 ratio in *b*-type to *a*-type character.

D. Modeling the HOONO ν_1 Band. With these ab initio predictions, we modeled the expected rotational contours of the ν_1 band of *cis-cis* HOONO using the asymmetric rotor program ASYROT.⁷⁹ We made the rigid rotor approximation and used the same set of rotational constants for the *cis-cis* isomer (from the ab initio calculations) for both the upper and lower states, because we had no information on the change in geometry upon excitation of the $\nu_1 = 0 \rightarrow 1$ transition.

To compare the simulation with the 3306 cm^{-1} spectrum, we first had to subtract the contributions to the observed spectrum from the adjacent $2\nu_2$ bands of *cis*-HONO and *trans*-HONO, whose wings overlapped with the HOONO band. We modeled the rotational contours of the HONO bands,⁷⁹ using asymmetric rotor constants from high-resolution FTIR data.^{80,81}

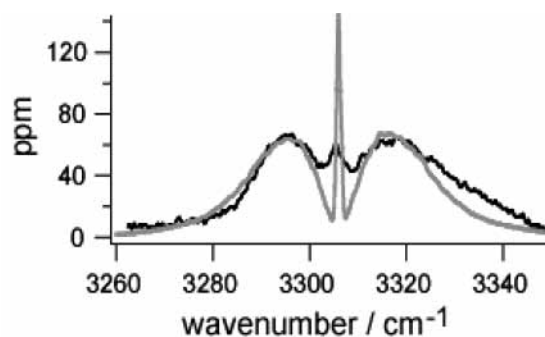


Figure 5. Comparison of the experimental *cis-cis* HOONO ν_1 spectrum (with background and HONO $2\nu_2$ overtone bands subtracted) and a simulation (grey line) computed as a hybrid band with ratio $I_b/I_a = 2.3$ using ab initio rotational constants.

The modeled spectra matched reference spectra of HONO taken with this apparatus reasonably well. Figure 4 shows the fit of the modeled *cis*- and *trans*-HONO spectra.

In Figure 5, a scan of HOONO from 3245 to 3380 cm^{-1} taken at 0.2 cm^{-1} steps with the background and HONO contributions subtracted, is compared with the simulation based on the ab initio predictions and a I_b/I_a ratio of 2.3. The simulation is in good agreement with the observed spectrum. We varied the orientation of the transition dipole moment, but found that the spectrum could best be fit as a hybrid band of *a*-type and *b*-type transitions. A range of I_b/I_a values from 1 to 4 was acceptable, because the Q branch blended with the P and R branches, making it difficult to estimate the relative area of the Q branch. This range is consistent with the predicted ratio $I_b/I_a = 2.3:1$.

The observed HOONO spectrum has a broader Q branch which partially blends with the P and R branches and a broader R branch shaded to the blue. We attribute these discrepancies to (a) the use of identical upper and lower rotational constants, and (b) contributions from sequence bands, progressions, or hot bands of the low-frequency modes of the loosely bound ring structure. The Q branch, in particular, is not expected to be as sharp as predicted by the simple rigid-rotor approximation, but the relative intensity roughly does match the observed spectrum. Blue-shading of the spectrum may arise from intramolecular hydrogen bond tightening in the vibrationally excited state, leading to a more compact geometry and slightly larger upper state rotational constants. Such effects are well documented in hydrogen-bonded complexes. More important, however, will be the significant population in excited low-frequency vibrations. Our CCSD(T)/cc-pVTZ calculations predict three low-frequency modes—the *a''* HOON torsion, the *a''* ONOO torsion, and the *a'* NOO bend—at 383 cm^{-1} , 523 cm^{-1} , and 419 cm^{-1} , respectively. The OH torsion, in particular, will break the hydrogen bond and lead to substantially blue-shifted sequence bands.

Figure 4 shows the total fit of the summed HOONO and HONO contributions. Together, the three spectra account for most of the extinction in this region. If we assume that the remaining area can be attributed to HOONO sequence bands, we estimate the total HOONO integrated signal to be 14% greater than the area of the simulated spectrum alone. This additional area is primarily on the blue side, and is comparable to the expected contribution from torsionally excited modes. The additional absorption is thus most likely to be due to sequence bands.

E. Lifetime of HOONO. The detection of *cis-cis* HOONO was performed at room temperature downstream of the reaction region, well after reactions 1 and 2 were complete. Typical total

residence times were 200 ms; thus, any measurements of absorption by HOONO could only be correlated with HOONO concentration if the HOONO lifetime relative to decomposition at room temperature was significantly longer than 200 ms. We varied the residence time from 75 to 250 ms, but we saw no change in the absolute HOONO signal. If the lifetime of HOONO were $\tau = 200$ ms, we would have expected to see the HOONO signal change by a factor of 1.7 over this time range. Given an uncertainty in HOONO detection of 20%, we estimate that the lifetime must be $\tau > 500$ ms. Thus, we conclude that the HOONO concentration was not changing appreciably over the time scale of this experiment. These conclusions apply as well to mechanisms which might increase *cis-cis* HOONO population, e.g., isomerization of *trans-perp* HOONO to *cis-cis* HOONO during transit through the detection cell.

F. Characterization of Wall Losses. In addition to decomposition of *cis-cis*-HOONO, loss of either HOONO or HONO₂ on the walls of the cell would result in erroneous product yields. While we saw no significant HOONO loss processes at room temperature, we found small decreases in the HOONO and HONO₂ signal at temperatures between 260 and 280 K.

These losses, on the order of 10–20%, increased when a cooled extension was inserted, indicating that some wall deposition or reaction was occurring. At $T < 250$ K, substantial reversible deposition of HONO₂ occurred. Our results are in qualitative agreement with unpublished wall loss measurements for HOONO by Fry et al.⁸² We therefore limited our HOONO/HONO₂ branching measurements to temperatures above 270 K.

G. Integrated Absorbances of HOONO and HONO₂. Integrated absorbances of the ν_1 bands of both HOONO and HONO₂ were obtained by scanning the respective ν_1 bands several times, recording discharge-on and discharge-off spectra, taking the difference, and then summing the runs. The bands were then numerically integrated.

We typically scanned each HOONO spectrum over the range 3250 to 3375 cm^{-1} and fit the observed bands to a composite spectrum from the simulated *cis*-HONO *trans*-HONO, and *cis-cis* HOONO spectra, allowing contributions from each of the three molecules to vary independently. We then used the area under the HOONO curve, scaled to account for the total area of the observed HOONO spectrum in Figure 5, to obtain the integrated HOONO absorption

$$I_{\text{HOONO}} = L_s \int \alpha_{\text{HOONO}}(\nu) d\nu$$

We obtained integrated absorption intensities of nitric acid in a similar manner, scanning from 3490 to 3610 cm^{-1} . It was essential to subtract the discharge-off spectrum, as there was a significant amount of background HONO₂ accounting for up to 10–15% of the total nitric acid. The spectrum in this region consisted of the main fundamental ν_1 band of HONO₂, as well as a sequence band, and a small contribution from the ν_1 band of *trans*-HONO. To obtain the integrated absorption of HONO₂, a spectrum of pure bubbled HONO₂ of comparable intensity was fit to our spectra and then integrated. We verified that a combination of the bubbled HONO₂ and *trans*-HONO spectra adequately represented the observed spectral intensity in the region from 3490 to 3610 cm^{-1} .

The magnitude of the HONO₂ absorption raised concerns about possible systematic errors in the observed absorbance of the nitric acid spectrum. The peak HONO₂ absorption was $\alpha l \leq 0.002$ (at the Q branch), so that the spectrum did not appear to be saturated in a single pass. However, the HONO₂ spectra

were typically 50 times stronger than the HOONO spectra, and the resulting ringdown lifetime was an order of magnitude greater than the baseline ringdown time.

We therefore performed a series of experiments to examine the ringdown spectra of varying concentrations of HONO₂ from nitric acid mixed with Ar at 20 Torr. We found that the overall absorption band shape varied slightly with HONO₂ concentration. The band shape agreed well with modeled contours at low HONO₂ concentrations, but the P and Q branches appeared to decrease relative to the R branch at higher concentrations. Hence, some apparent nonlinearity in the absorption measurement appeared to be occurring.

The use of a low-resolution laser for cavity ringdown spectroscopy may introduce a potential error. As Zalicki and Zare⁶³ and Yalin and Zare⁸³ have shown, when the ratio of the absorption bandwidth to laser bandwidth is $V < 100$, the ringdown decay becomes nonexponential. At a given frequency, the laser samples some ensemble of narrow, partially resolved, and unresolved absorption features within the line width of 1 cm^{-1} . The total signal is the sum of frequency components ν_i , each contributing a signal with a decay time τ_i determined by the absorption $\alpha(\nu_i)$. The resulting ringdown signal in general will not be a single exponential. In a typical experiment, the total signal is then a fit to an effective overall decay, which underestimates the absorption. As the spectrum becomes increasingly congested, however, narrow features blend and the lines can become unresolvable. If the spectral features become much wider than the laser line width, the cavity ringdown spectrum should approach the convoluted low-resolution spectrum. Alternatively, Zalicki and Zare also showed that if the absorbance was weak compared to cavity loss, then the integrated intensity—though not the individual line shape—was correct to within 10%. Furthermore, as Hodges et al.⁸⁴ have shown, the spectral profile of YAG-pumped dye lasers is not Gaussian, but composed of fluctuating lines at the cavity mode spacing, which tends to ameliorate the underestimation.

The ν_1 band of nitric acid is known to have resolvable Doppler-broadened features, but it is highly congested and has not yet been assigned.⁸⁵ In the current experiment, the underlying HONO₂ spectrum typically consisted of a set of narrow lines with a Voigt profile from the convoluted Doppler and pressure-broadened line shapes. The ringdown spectra that we recorded possessed partially resolved features.

We therefore determined the deviation of the integrated absorbance measured by CRDS in our branching ratio experiments in a separate experiment, in which we measured the absolute integrated intensity of the ν_1 band as a function of HONO₂ concentration. N₂ was bubbled through nitric acid and its concentration measured by a 185 nm Hg lamp before dilution with N₂ and passage into the CRDS cell. A detailed description of this experiment, as well as simulations of ringdown spectra based on high-resolution spectra of HONO₂, will be reported elsewhere.⁸⁶ The true integrated intensity was measured at low HONO₂ concentrations, and compared to HONO₂ absorbance measured at the HONO₂ concentrations obtained in the discharge. From these measurements, we concluded that the current branching ratio experiments underestimated the absolute infrared intensity of the ν_1 band of HONO₂ by 17%.

In contrast, it is likely that the integrated absorbance of the ν_1 band of *cis-cis* HOONO is correct. First, the HOONO absorption is approximately 10 times less than the mirror loss, thus meeting the criterion of Zalicki and Zare for weak absorption. Second, we expect that this band is highly congested at room temperature. The observed spectrum is relatively

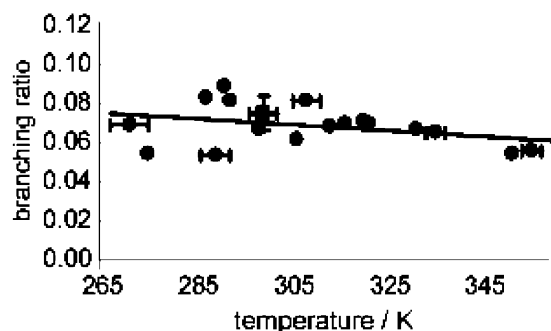


Figure 6. Observed ratio of *cis-cis* HOONO to HONO₂ products in the reaction OH + NO₂ as a function of temperature, at 20 Torr (See Table 1 for gas composition). The experimentally measured ratios of the integrated ν_1 absorption bands were scaled by the computed ratio of absorption cross sections. The abscissa uncertainties are the 2σ random errors associated with temperature variations. One point (square) was the mean of measurements on several days.

smooth, with no detectable sub-structure. *cis-cis* HOONO is predicted to be an asymmetric rotor with Ray's parameter $\kappa = -0.71$, and the ν_1 band is predicted to be an *a-b* hybrid. Since the OH bond participates in an intramolecular hydrogen-bond, the rotation-vibration interactions cannot be neglected. Finally, several low-frequency modes, including two torsional modes, will have significant populations that will increase the spectral congestion. HOONO has three modes with harmonic frequencies under 500 cm⁻¹, whereas the lowest frequency harmonic mode in nitric acid is slightly below 500 cm⁻¹. Under these conditions, the "low-resolution" cavity ringdown spectrum of HOONO is more likely to both reproduce the overall shape of the band and yield the correct integrated absorbance.

H. Determination of the *cis-cis* HOONO/HONO₂ Branching Ratio. We recorded the integrated absorptions of *cis-cis* HOONO and HONO₂ at a variety of reaction temperatures at 20 Torr, and pressures from 4 to 40 Torr at 300 K. While the temperature of the reaction region was varied, all spectroscopic measurements were performed with the detection cell held at room temperature. Thus, the band contours were independent of reaction conditions, and wall losses in the detection region were negligible. Furthermore, as we demonstrated above, the HOONO lifetimes were sufficiently long that HOONO extinction coefficients were unaffected by the residence time. Thus, the column density in the spectroscopy should accurately reflect the product yield. The HONO₂ integrated absorption was scaled by 1.17 to correct for the error due to cavity ringdown measurement. The HONO₂ concentrations did not vary considerably in the data sets, and we found that use of the same correction factor did not introduce a significant error.

HOONO-to-HONO₂ branching ratios were then obtained by weighting the observed absorption intensity ratios by 2.87, the ratio of the *ab* initio absorption band strengths computed at the CCSD(T)/cc-pVTZ level.

Measurements of the branching ratio were performed at room temperature and 20 Torr on several different days. The variations, most likely arising from source conditions, were ± 0.025 . Averaging over these results, we determined the 300 K branching ratio to be $k_2(c-c)/k_1 = 0.075 \pm 0.009$, where the quoted uncertainty is the 2σ standard error of the mean.

Figure 6 shows our measurements of the branching ratio of *cis-cis* HOONO to HONO₂ over the temperature range 270 K to 360 K at 20 Torr total pressure. The *cis-cis* HOONO yield appeared to be almost independent of temperature; the slope of the temperature profile was $-1.7(\pm 2.4) \times 10^{-4}/\text{K}$ (2σ). The temperatures plotted were from the thermocouple in the reaction

region. Typical uncertainties in T due to radial variations of up to ± 3 K are shown, but the weak temperature dependence leads to little resulting uncertainty in the yield of HOONO.

I. Uncertainties. The scatter in the data shown in Figures 6 is due to random errors that include the effects of (a) noise in the collected spectra; (b) uncertainty in the estimation of the integrated band intensities from fitting the band of HOONO, HONO, and HONO₂; and (c) variations in conditions of the reaction region. As discussed above, there is a temperature uncertainty of up to ± 3 °C. Our uncertainty in the pressure is 1%. However, the branching ratio depends only weakly on temperature, and thus the error propagated into the branching ratios is negligible.

In addition to the random experimental errors, there are several uncertainties. The largest uncertainty arises from the calculated ratio of absorption cross sections, which we estimate is on the order of 20% on the basis of the difference between the computed and observed ratios of ν_1 band strengths of HONO₂ and HO₂NO₂. A second uncertainty lies in our correction of the integrated HONO₂ absorbance due to the use of the cavity ringdown method. We estimate the error in this correction to be on the order of 10%. Finally, if we conservatively assume that there is resolvable structure in the HOONO spectrum, then the weak absorbance condition of Zalicki and Zare⁶³ leads to 10% uncertainty in the HOONO intensity. Adding these uncertainties in quadrature to the 2σ random error of 0.009(12%), we find that the branching ratio for *cis-cis* HOONO vs HONO₂ products is

$$k_2(c-c)/k_1 = 0.075 \pm 0.020(2\sigma)$$

at room temperature and 20 Torr of the He/Ar/N₂ mixture.

J. Search for ν_1 Band of *trans-perp* HOONO. Our spectroscopic measurements in the 3500–3650 cm⁻¹ range can also provide a coarse upper bound on the *trans-perp* HOONO concentration. From the matrix isolation studies, the ν_1 band of *trans-perp* HOONO is near 3550 cm⁻¹, under the OH stretch of nitric acid. Our calculations predict that this band is 50% weaker than the ν_1 band of HONO₂. This band is likely to be a *c*-type transition, because the OH bond lies nearly perpendicular to the plane of the heavy atoms. Thus, the *trans-perp* OH stretch band should exhibit a sharp Q branch that would be readily observed within the HONO₂ absorption (unless it lies directly under the HONO₂ Q branch). From the absence of any anomalous features or other large discrepancies in the vicinity of the HONO₂ spectrum, we infer that the branching fraction of the *trans-perp* isomer is $< 20\%$.

K. Statistical Calculation of the *trans-perp* → *cis-cis* Isomerization Rate. We simulated the isomerization process, assuming that these isomers are separated by the transition state with geometry corresponding to the *perp-perp* saddle point computed at the MP2/6-311G(3df,2p) level by McGrath and Rowland.³⁷ This transition state has an energy of 11.4 kcal/mol relative to the *cis-cis* isomer at the MP2 level, including zero-point energy. For consistency in the density of states calculations, we used their MP2 vibrational frequencies and geometries for all three conformers, *cis-cis*, *trans-perp*, and *perp-perp*, but used our best CCSD(T) estimate for the *trans-perp* energy relative to the energy of *cis-cis* HOONO.

We found that the nascent HOONO molecules establish a rapid ($k > 10^{10}$ s⁻¹) microcanonical equilibrium between the two isomers as they are slowly ($k_{\text{collision}} \approx 10^8$ s⁻¹ at 13 Torr) driven down the internal energy ladder by collisions. Because of its smaller moments of inertia (centrifugal barrier effect) and larger degree of "floppiness", the *trans-perp* isomer has a density

of states comparable to that of the *cis-cis* isomer at energies above the isomerization barrier. Thus, a large fraction of excited HOONO molecules ($\approx 30\%$) initially fall into the *trans-perp* well. Once the internal energy of the HOONO molecules is below the isomerization barrier, vibrational energy is equilibrated with translations on a microsecond time scale under our experimental conditions. This initial yield of vibrationally thermalized HOONO isomers depends critically on the relative density of states of the two isomers just above the top of the barrier at low pressures. At higher pressures, collisional relaxation may become competitive with isomerization rates.

After the HOONO molecules relax to an initial *c-c/t-p* population, the *trans-perp* HOONO isomerize on a slower time scale of 0.1–10 ms (over the temperature range of interest), and the *c-c/t-p* populations approach thermal equilibrium. Using the 11.4 kcal/mol barrier for the *cis-cis* to *trans-perp* isomerization and the 3.4 kcal/mol energy separation between the isomers (both zero-point corrected) and assuming a doubly degenerate reaction path, we find isomerization lifetimes of 0.35 ms and 1.5 ms at 300 and 260 K, respectively. Even with a zero-point corrected *perp-perp* energy of 13.5 kcal/mol and assuming a single reaction path, the lifetime for *trans-perp* relative to isomerization was 18 ms at 300 K. Thus, *trans-perp* HOONO is expected to convert to the lowest energy conformer rapidly relative to the characteristic residence time of products in the flow cell (200 ms). Therefore, the initial yields of the HOONO isomers from $\text{OH} + \text{NO}_2 + \text{M}$, which are difficult to compute, will not be detected in our experiment.

Discussion

A. Total HOONO Yield. The current experiment measures the branching ratio for the formation of one HOONO isomer, the lowest energy *cis-cis* conformation, relative to formation of nitric acid. These results thus put a rigorous lower bound on the total HOONO yield, as other isomers of HOONO may have been formed but were undetected.

Two other HOONO isomers have been predicted, both with OH stretch bands that would overlap with that of nitric acid. Recent calculations conclude that the second lowest energy form, the *cis-perp* isomer, is an inflection point, not a minimum, along the torsional potential about the peroxy bond. This result is consistent with the absence of spectroscopic evidence for the *cis-perp* isomer in matrixes. We consider first the possible formation of the *trans-perp* isomer, whose energy is predicted to be $\Delta E_0 = +3.4$ kcal/mol relative to the *cis-cis* isomer, and then briefly discuss possible problems arising from torsionally excited *cis-cis* HOONO.

Any *trans-perp* HOONO products formed by reaction 2 would most likely isomerize to the *cis-cis* form, because the computed isomerization barrier (11.4 kcal/mol) is predicted to be much lower than the dissociation threshold (18.3 kcal/mol). The absence of any change in the *cis-cis* HOONO abundance with residence time is consistent with one of three hypotheses: (a) the *trans-perp* form equilibrates rapidly with *cis-cis*, prior to entry into the detection volume, (b) no *trans-perp* isomer is formed, or (c) the *trans-perp* isomer isomerizes very slowly compared to the residence time.

Our RRKM calculations with MultiWell support the first scenario, that *trans-perp* isomers rearrange rapidly and are in equilibrium with *cis-cis* HOONO in the detection volume. We calculate that the rate is $>10^3 \text{ s}^{-1}$ at room temperature, and is sufficiently fast, even if the zero-point corrected barrier is as high as 13.5 kcal/mol. Furthermore, the calculations assume that the *perp-perp* conformation is the only transition state, but other pathways may exist.

If isomerization of the *trans-perp* isomer to the *cis-cis* form is indeed rapid compared to the residence time, we can calculate the total yield of HOONO in reactions 1 and 2. From the energetics and frequencies predicted by our CCSD(T) calculations in Table 2, we find $K_{\text{eq}}(300 \text{ K}) = [\textit{trans-perp} \text{ HOONO}] / [\textit{cis-cis} \text{ HOONO}] = 0.005$. Thus, at equilibrium only a negligible fraction the HOONO will be in the *trans-perp* conformation. Even at the highest temperature reported here, the *trans-perp* isomer contributes less than 1.5% to the total population at equilibrium.

While the *cis-perp* isomer is not bound in CCSD(T) calculations,⁴⁶ its geometry may be accessed by low ($\nu_9 \geq 1$) torsional levels of the NOOH torsion. This torsional motion breaks the intramolecular hydrogen bond, and the large predicted change in the OH vibrational frequency with torsional angle suggests that the sequence bands involving ν_9 will be significantly blue-shifted. The excited torsional modes—which may account for 10–30% of the *cis-cis* HOONO population—may be the cause of the intensity on the blue side of the spectrum in Figure 5. However, we cannot rule out that some of the sequence bands cannot be detected within the 3306 cm^{-1} band. Thus, the observed yield obtained from the integrated absorbance of the 3306 cm^{-1} band may underestimate the HOONO yield, though by at most 10–30%. Only detailed modeling of the torsional levels with accurate ab initio potentials, or direct observation of the sequence bands, will allow us to settle this question.

We conclude that the observed *cis-cis* yield accounts for essentially all HOONO formed, i.e., that $k_2/k_1 = 0.075$ at 20 Torr and 300 K for the reaction $\text{OH} + \text{NO}_2 + \text{M}$.

B. Comparison with Previous Measurements. Both Burkholder et al. and Dransfield et al. were unable to detect HOONO by FT-IR spectroscopy. Burkholder et al. placed an upper limit of 5% on the HOONO yield,⁵¹ but assumed equal ν_1 cross sections with nitric acid; if our CCSD(T) intensities are used, this upper limit becomes 14%. Our yield also agrees with their HONO₂ yield of $75\% + 25\% / -10\%$ determined from mass balance. Dransfield et al. scanned a wider spectral range under a variety of conditions;⁵² they reported an estimate from the mass balance of the HOONO yield to be $0.0 \pm 4.1\%$ at 298 K and 60 Torr. Curiously, their published spectra end at 3200 cm^{-1} . Our experiments suggest that the concentrations of HOONO in their experiments were near the limits of the reported sensitivities. In addition, their measurements were performed at higher pressures, where HOONO yields may be higher but collisional dissociation will be faster.

Nizkorodov and Wennberg⁵⁴ reported a HOONO yield of $5 \pm 3\%$ at 253 K and 20 Torr. These measurements were performed under conditions similar to our own, but were based on detection of HOONO alone. Their uncertainty arose because the authors assumed that the HOONO overtone cross-sections were the same as that of HO₂NO₂, and that HOONO dissociated with unit quantum efficiency. There was also some uncertainty in the assignment of the peaks, as no rotational structure was resolved. They assigned peaks in a cluster of bands to the ν_1 overtones of the *cis-cis*, *cis-perp*, and *trans-perp* isomers; more recent results⁸² suggest that these are all combination bands, hot bands, and/or Fermi resonances of the *cis-cis* isomer. Given these assumptions, their yield is in quite good agreement with the current findings.

Donahue et al. estimated a low-pressure branching ratio of 0.17 from their isotope exchange measurements, but this estimate is rather weakly constrained by data.⁵³ Furthermore, since the high-pressure data had not been corrected, they

consequently overestimated k_2^∞ for HOONO, and this discrepancy will have had an impact on their low-pressure extrapolation.

Hippler et al. measured a high-pressure branching ratio in He of 0.33 at 430 K, which fell to 0.15 at 5 bar He.³² Since HOONO yields are predicted to be further reduced at lower pressures, these results are qualitatively consistent with the present data. They did not estimate rate constants in the low-pressure limit, so a direct comparison would require the pressure and temperature dependences.

Comparisons with the earlier modeling results are more tenuous, since the calculations are sensitive to input parameters. Troe's calculation,³⁴ for example, may have underestimated the HOONO yield by using too low a binding energy. Matheu and Green³³ attempted to model both forms of HOONO; they find the low-pressure ratio $k_2/k_1 \geq 0.02$. Our results are in good agreement with Golden and Smith,³¹ who predict $k_2/k_1 = 0.1$ in the low-pressure limit; this rate was arrived at by multiplying the computed *cis-cis* rate by three to account for the three isomers. Our results are in satisfactory agreement with recent work by Golden, Barker, and Lohr,³⁵ the most detailed master equation modeling calculations to date. From their fitted parameters, we compute a branching ratio of 0.057 at 13 Torr of N₂.

C. Atmospheric Implications. The low-pressure branching ratios reported here provide an important constraint on the low-pressure k_2^0/k_1^0 branching ratio for models, and are a lower limit on HOONO yields in the atmosphere. The gas mixture, He:Ar:N₂, at 20 Torr is equivalent to approximately 13 Torr of dry air after collision efficiencies β of the gases are taken into account. The HOONO yield at pressures relevant to the atmosphere will be higher than the yields reported here. Doubling or tripling of this yield could significantly affect atmospheric models. However, while calculations predict that HOONO yields will increase to high levels at atmospheric pressures, one cannot yet predict with full confidence the true falloff curves for HOONO.

Laboratory measurements of reaction 1 generally detect only loss of reactants (typically OH under pseudo-first-order conditions); however, the reported rate coefficient may be either k_1 or $(k_1 + k_2)$, depending on whether HOONO has collisionally dissociated on the time scale of the measurements. The implications of the current branching ratio will then depend on the specific experiment.

Conclusions

This work establishes the formation of HOONO as a minor product channel in the reaction OH + NO₂, consistent with recent experimental observations by Donahue,⁵³ Nizkorodov,⁵⁴ and Hippler.³² Our experiments unambiguously identify HOONO by mid-infrared spectroscopy and provide a direct measure of the branching ratio in the temperature range 270–360 K. We have exploited the sensitive technique of cavity ringdown spectroscopy to detect the absorption spectrum of the ν_1 band of *cis-cis* HOONO, centered at 3306 cm⁻¹. The rotational envelope, obtained after subtracting the contributions from the neighboring HONO ν_2 overtones, agrees well with the simulated spectrum derived from high-level *ab initio* calculations. We determined the HOONO:HONO₂ product branching ratio k_2/k_1 for the OH + NO₂ + M reaction as a function of reaction pressure and temperature from measurements of the integrated absorption of the ν_1 bands of *cis-cis* HOONO and HONO₂. These measurements were converted to branching ratios using the ratio of calculated integrated intensities of HOONO and HONO₂ bands at the CCSD(T)/cc-pVTZ level. The use of

computed cross sections leads to the largest contribution to the uncertainty in the branching ratio. We find a branching ratio of $k_2/k_1 = 0.075 \pm 0.020(2\sigma)$ at room temperature and 20 Torr (He/Ar/N₂ buffer gas). This branching ratio decreases weakly with increasing temperature over the range 270–360 K. Although a second isomer of HOONO may be formed, statistical rate calculations show that the higher energy conformer, *trans-perp* HOONO, is expected to isomerize rapidly on the time scale of the current experiment, and will be in its equilibrium abundance, computed to be < 1.5% of the total population. A small fraction of ν_9 (HOON-torsion) excited molecules that access the *cis-perp* geometry may have sequence bands that are significantly blue-shifted and lie under the HONO₂ band, and thus be undetected. Thus, the current measurements provide a firm lower bound for and may be a good estimate of the total yield of HOONO.

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